# Hilbert Norms For Graded Algebras

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### Abstract

This paper presents a solution to a problem from superanalysis about the existence of Hilbert-Banach superalgebras. Two main results are derived: 1) There exist Hilbert norms on some graded algebras (infinite-dimensional superalgebras included) with respect to which the multiplication is continuous. 2) Such norms cannot be chosen to be submultiplicative and equal to one on the unit of the algebra.

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## 1 Introduction

The type of norms investigated in this article are generalizations of norms used for the symmetric tensor algebra in the white noise analysis [7][11] or in the Malliavin calculus [20]. But now more general algebras are included, especially the algebra of antisymmetric tensors (Grassmann algebra) and  $\mathbb{Z}_2$ -graded algebras (superalgebras) related to supersymmetry and to quantum probability [15].

A locally convex commutative superalgebra is a  $\mathbb{Z}_2$ -graded locally convex space  $\mathcal{E} = \mathcal{E}_0 \oplus \mathcal{E}_1$  equipped with an associative continuous multiplication having the following property: for any  $a, b \in \mathcal{E}_0 \cup \mathcal{E}_1$ ,  $ab \neq 0$  the product satisfies  $ab = (-1)^{p(a)p(b)}ba$  with the parity function p, which is defined on  $(\mathcal{E}_0 \cup \mathcal{E}_1) \setminus \{0\}$  with  $p(\mathcal{E}_0 \setminus \{0\}) = 0$ ,  $p(\mathcal{E}_1 \setminus \{0\}) = 1$ , and p(ab) = |p(a) - p(b)|. Typical examples are Grassmann algebras with finite or countable sets of generators. In superanalysis one considers modules over (commutative) superalgebras [16][8][5][19][17][4][18][10].^3 It is quite easy to define an infinite-dimensional Grassmann algebra with a non-Hilbertian norm [16]. But for a long time it was unknown whether the topology of a locally convex superalgebra - including the Grassmann algebra. The paper gives a complete solution to these problems. Our theorems imply a positive answer to the first question and a negative answer to the second question.

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<sup>&</sup>lt;sup>3</sup>In the pioneering works of Martin [14] and of Berezin [3] the Grassmann algebra itself has been used instead of these modules.

#### $\mathbf{2}$ General considerations

Let  $\mathcal{A}$  be an algebra over the field  $\mathbb{K} = \mathbb{R}$  or  $\mathbb{C}$  with unit  $e_0$ . The product is denoted by  $a, b \in \mathcal{A} \to ab \in \mathcal{A}$ . We assume that  $\mathcal{A}$  is provided with a positive definite inner product  $a, b \in \mathcal{A} \to (a \mid b) \in \mathbb{K}$ . The corresponding Hilbert norm  $||a|| = \sqrt{(a \mid a)} \ge 0$  is normalized at the unit  $||e_0|| = 1$ . We are interested in such norms which allow a uniform estimate for the product of the algebra

$$\|ab\| \le \gamma \|a\| \|b\| \tag{1}$$

with a constant  $\gamma \geq 1$ . In this section we prove under rather general conditions that this constant has the lower limit  $\gamma \ge \sqrt{\frac{4}{3}}$ .

**Theorem 1** Let  $\mathcal{A}$  be an algebra over the field  $\mathbb{K} = \mathbb{R}$  or  $\mathbb{C}$  with dimension dim  $\mathcal{A} \geq 2$ . If this algebra satisfies the properties

i)  $\mathcal{A}$  is provided with a Hilbert inner product (. |.) normalized at the unit  $e_0$ ,  $||e_0||^2 = 1$  $(e_0 \mid e_0) = 1,$ 

ii) there exists at least one element  $f \in \mathcal{A}, f \neq 0$ , such that  $e_0, f$  and  $f^2 = ff$  satisfy  $(e_0 \mid f) = (f \mid f^2) = 0 \text{ and } (e_0 \mid f^2) \ge 0,$ 

then the norm estimate  $||ab|| \leq \gamma ||a|| ||b||$  is not valid for some  $a, b \in \mathcal{A}$ , if  $\gamma < \sqrt{\frac{4}{3}}$ .

**Proof** Since  $f \neq 0$  we can normalize this element and assume ||f|| = 1. Take  $a = e_0 + \lambda f$ 

with  $\lambda \in \mathbb{R}$ . Then  $a^2 = e_0 + 2\lambda f + \lambda^2 f^2$  and  $\|a^2\|^2 = 1 + 2\lambda^2 (e_0 | f^2) + 4\lambda^2 + \lambda^4 \|f^2\|^2 \ge 1 + 4\lambda^2$ . On the other hand  $\|a\|^2 = 1 + \lambda^2$ , and  $\|a^2\|^2 \le \gamma^2 \|a\|^4$  implies  $1 + 4\lambda^2 \le \gamma^2 (1 + \lambda^2)^2$ . But this inequality is true for all  $\lambda \ge 0$ only if  $\gamma^2 \ge \sup_{\lambda \ge 0} (1 + 4\lambda^2)(1 + \lambda^2)^{-2} = \frac{4}{3}$ .

This Theorem obviously applies to the tensor algebra  $\mathcal{T} = \bigoplus_{n=0}^{\infty} \mathcal{T}_n$ , where  $\mathcal{T}_n$  is the subspace of tensors of degree n, and the norm is defined in the standard way as

$$||f||^{2} = \sum_{n=0}^{\infty} w_{n} ||f_{n}||_{n}^{2} \text{ if } f = \sum_{n=0}^{\infty} f_{n}, \ f_{n} \in \mathcal{T}_{n}$$
(2)

with arbitrary positive weights  $w_n > 0, n \in \mathbb{N}$  and  $w_0 = 1$ . In that case we can simply choose an element  $f \in \mathcal{T}_1, f \neq 0$ , to satisfy the assumptions with  $(e_0 \mid f \otimes f) = 0$ .

Theorem 1 can also be applied to a large class of algebras  $\mathcal{A}$  which can be derived from the tensor algebra  $\mathcal{T}$  by the following modifications of the product.

- 1. The product is generated by  $f, g \in \mathcal{A}_1 = \mathcal{T}_1 \to f \circ g := f \otimes g + (-1)^{\chi} g \otimes f$  where  $\chi = 0, 1 \mod 2$  is a parity factor.
- 2. The product is generated by  $f, g \in \mathcal{A}_1 = \mathcal{T}_1 \rightarrow$  $f \circ g := f \otimes g + (-1)^{\chi} g \otimes f + \omega(f,g) e_0$ . Here  $\chi$  is again a parity factor and  $\omega(.,.)$ :  $\mathcal{A}_1 \times \mathcal{A}_1 \to \mathbb{K}$  is a bilinear pairing.

The first class of algebras includes the algebra of symmetric tensors, the algebra of antisymmetric tensors (Grassmann algebra), and tensor products of these algebras, including the  $\mathbb{Z}_2$ -graded algebras (superalgebras) used in quantum field theory. The assumptions of the Theorem 1 are satisfied for any non-vanishing element  $f \in \mathcal{A}_1 = \mathcal{T}_1$ .

The second class includes the Clifford product, the (symmetric) Wiener product, the antisymmetric Wiener product (with antisymmetric  $\omega$ ) and Le Jan's supersymmetric Wiener-Grassmann product [9][13][15]. In these cases the assumptions of Theorem 1 are satisfied if there exists a non-vanishing  $f \in \mathcal{A}_1$  with  $\omega(f, f) \geq 0$ . Such a vector can always be found

if the algebra is complex, or

if the algebra is real and  $\omega$  is not negative definite.

The last constraint is satisfied for the symmetric Wiener product on real spaces, and for the real Clifford system in quantum field theory [2]. In both cases the form  $\omega$  is positive definite.

Moreover Theorem 1 is obviously true for any unital algebra  $\mathcal{A}$ , which has a nilpotent element f that is orthogonal to the unit element. If we only know that  $\mathcal{A}$  has at least one nilpotent element, we can derive the weaker

**Corollary 1** Let  $\mathcal{A}$  be an algebra which satisfies condition i) of Theorem 1. If this algebra has a nilpotent element f, then the norm estimate  $||ab|| \leq ||a|| ||b||$  is not valid for some  $a, b \in \mathcal{A}$ .

**Proof** We assume again ||f|| = 1. Then  $a = e_0 + \lambda f$  with  $\lambda \in \mathbb{R}$  and  $a^2 = (e_0 + \lambda f)^2 = e_0 + 2\lambda f$  have the norms  $||a||^2 = 1 + 2\lambda \operatorname{Re}(e_0, f) + \lambda^2$  and  $||a^2||^2 = 1 + 4\lambda \operatorname{Re}(e_0, f) + 4\lambda^2$ . If  $\operatorname{Re}(e_0, f) = 0$  we can apply the arguments given in the proof for Theorem 1. If  $\operatorname{Re}(e_0, f) = \gamma \neq 0$ , then we chose  $\lambda = -2\gamma$ , and  $||a^2||^2 = 1 + 8\gamma^2 \leq 1 = ||a||^4$  is a contradiction.

## 3 Norm estimates for $\mathbb{Z}$ -graded algebras

In this section we present Hilbert norm estimates for rather general  $\mathbb{Z}$ -graded algebras  $\mathcal{A}$  over the field  $\mathbb{K} = \mathbb{R}$  or  $\mathbb{C}$ . We assume the following structure of  $\mathcal{A}$ .

- 1. The algebra is the direct sum  $\mathcal{A} = \bigoplus_{n=0}^{\infty} \mathcal{A}_n$  of orthogonal spaces  $\mathcal{A}_n$ . Thereby  $\mathcal{A}_0$  is the one dimensional space  $\mathbb{K}$  spanned by the unit  $e_0$  of the algebra. The product  $a \circ b$  maps  $\mathcal{A}_p \times \mathcal{A}_q$  into  $\mathcal{A}_{p+q}$  for all  $p, q \in \{0, 1, ...\}$ .
- 2. The spaces  $\mathcal{A}_n$  are provided with Hilbert norms  $\|.\|_n$ ,  $n = 0, 1, \dots$  The unit has norm  $\|e_0\|_0 = 1$ . The product of two homogeneous elements  $a_p \in \mathcal{A}_p$  and  $b_q \in \mathcal{A}_q$  satisfies

$$\|a_{p} \circ b_{q}\|_{p+q} \le \|a_{p}\|_{p} \|b_{q}\|_{q}$$
(3)

if  $a_p \in \mathcal{A}_p$  and  $b_q \in \mathcal{A}_q$ .

3. The algebra is provided with a family of Hilbert norms

$$\|a\|_{(\sigma)}^{2} = \sum_{n=0}^{\infty} w_{n}(\sigma) \|a_{n}\|_{n}^{2} \text{ if } a = \sum_{n=0}^{\infty} a_{n}, \ a_{n} \in \mathcal{A}_{n}$$
(4)

with  $\sigma \in \mathbb{R}$ . The factors  $w_n(\sigma), n = 0, 1, ...,$  are positive weights with the normalization  $w_0(\sigma) = 1$  for all  $\sigma \in \mathbb{R}$ . The weights satisfy the inequalities  $w_n(\sigma) \leq w_n(\tau)$ for all  $n \in \mathbb{N}$  if  $\sigma \leq \tau$ .

An immediate consequence of these assumptions is  $\|a\|_{(\sigma)} \leq \|a\|_{(\tau)}$  for all  $a \in \mathcal{A}$  if  $\sigma \leq \tau$ . A simple example of such an algebra  $\mathcal{A}$  is the tensor algebra  $\mathcal{T}$ . Its standard norm satisfies (3) with weights  $w_n = 1$  for all n = 0, 1... More interesting examples are the algebras of symmetric tensors or of antisymmetric tensors. With the notation  $f \circ g$  for both the symmetric and the antisymmetric tensor product the estimate (3) is satisfied by the norms

$$\|f_1 \circ f_2 \circ \dots \circ f_n\|_n^2 = \begin{cases} (n!)^{-1} \operatorname{per}(f_\mu \mid f_\nu) \text{ for symmetric tensors,} \\ (n!)^{-1} \operatorname{det}(f_\mu \mid f_\nu) \text{ for antisymmetric tensors,} \end{cases}$$
(5)

but it is violated if the factor  $(n!)^{-1}$  is omitted. The standard norm<sup>4</sup> is defined without the factor  $(n!)^{-1}$ . In the notations used here it corresponds therefore to a norm (4) with a weight function  $w_n = n!$ .

**Theorem 2** If there exists a constant  $\delta(\sigma, \tau, \rho) > 0$  such that the weight functions satisfy the inequalities

$$(p+q-1)w_{p+q}(\rho) \le \delta(\sigma,\tau;\rho)w_p(\sigma)w_q(\tau) \quad if \ p,q \ge 1$$
(6)

for values of  $\sigma, \tau$  and  $\rho$  with  $\sigma \leq \rho$  and  $\tau \leq \rho$ , then the product of  $\mathcal{A}$  is estimated by

$$\|a \circ b\|_{(\rho)} \le \gamma \cdot \|a\|_{(\sigma)} \|b\|_{(\tau)} \tag{7}$$

where the constant  $\gamma$  is  $\gamma = \sqrt{3} \max(1, \delta(\sigma, \tau, \rho))$ .

**Proof** For  $a = a_0 + a_+$  and  $b = b_0 + b_+$  with  $a_0, b_0 \in \mathcal{A}_0 = \mathbb{K}$  and  $a_+ = \sum_{n=1}^{\infty} a_n$ ,  $b_+ = \sum_{n=1}^{\infty} b_n$  with  $a_n, b_n \in \mathcal{A}_n$ ,  $n \in \mathbb{N}$  the norm of  $a \circ b$  is calculated by

$$\begin{aligned} \|a \circ b\|_{(\rho)}^{2} &= \|a_{0}b_{0} + a_{0}b_{+} + a_{+}b_{0} + a_{+} \circ b_{+}\|_{(\rho)}^{2} \\ &\leq |a_{0}b_{0}|^{2} + 3\left(|a_{0}|^{2} \|b_{+}\|_{(\rho)}^{2} + \|a_{+}\|_{(\rho)}^{2} |b_{0}|^{2} + \|a_{+} \circ b_{+}\|_{(\rho)}^{2}\right) \\ &\leq |a_{0}b_{0}|^{2} + 3\left(|a_{0}|^{2} \|b_{+}\|_{(\rho)}^{2} + \|a_{+}\|_{(\rho)}^{2} |b_{0}|^{2} + \sum_{n \geq 1} w_{n}(\rho) \left\|\sum_{p+q=n}^{\prime} a_{p} \circ b_{q}\right\|_{n}^{2}\right) \end{aligned}$$

The symbol  $\sum'$  means summation with the constraint  $p \ge 1, q \ge 1$ . The sum  $\sum_{p+q=n, p\ge 1, q\ge 1} \ldots = \sum'_{p+q=n} \ldots$  has n-1 terms, hence  $\left\|\sum'_{p+q=n} a_p \circ b_q\right\|_n^2 \le (n-1) \sum'_{p+q=n} \|a_p \circ b_q\|_n^2 \stackrel{(3)}{\le} (n-1) \sum'_{p+q=n} \|a_p\|_p^2 \|b_q\|_q^2$ . If  $w_n(\rho)$  is chosen such that (6) is satisfied we obtain  $\sum_{n\ge 1} w_n(\rho) \left\|\sum'_{p+q=n} a_p \circ b_q\right\|_n^2 \le \delta \cdot \left(\sum_{p\ge 1} w_p(\sigma) \|a_p\|_p^2\right) \cdot \left(\sum_{q\ge 1} w_q(\tau) \|b_q\|_q^2\right)$ 

<sup>&</sup>lt;sup>4</sup>The "standard" inner product of the symmetric/antisymmetric tensor algebra is characterized by the following property. Let  $\mathcal{F}_i$ , i = 1, 2, be two orthogonal subspaces of the space  $\mathcal{A}_1$ . Denote by  $\mathcal{A}(\mathcal{F}_i)$ the subalgebra generated by elements  $f \in \mathcal{F}_i$ . Then  $(a_1 \circ a_2 \mid b_1 \circ b_2) = (a_1 \mid b_1) (a_2 \mid b_2)$  holds for all  $a_i \in \mathcal{A}(\mathcal{F}_i), i = 1, 2$ .

 $\leq \delta \|a_{+}\|_{(\sigma)}^{2} \|b_{+}\|_{(\tau)}^{2}.$  For  $\rho \leq \sigma, \tau$  we have in addition the inequalities  $\|a_{+}\|_{(\rho)}^{2} \leq \|a_{+}\|_{(\sigma)}^{2}$  and  $\|b_{+}\|_{(\rho)}^{2} \leq \|b_{+}\|_{(\tau)}^{2}$  such that finally

$$\|a \circ b\|_{(\rho)}^{2} \leq |a_{0}b_{0}|^{2} + 3\left(|a_{0}|^{2} \|b_{+}\|_{(\tau)}^{2} + \|a_{+}\|_{(\sigma)}^{2} |b_{0}|^{2} + \delta \|a_{+}\|_{(\sigma)}^{2} \|b_{+}\|_{(\tau)}^{2}\right)$$
  
$$\leq 3\gamma \|a\|_{(\sigma)}^{2} \|b\|_{(\tau)}^{2}.$$

where  $\gamma$  is  $\gamma = \max(1, \delta)$ .

As the first application of Theorem 2 we derive norms with respect to which the product of the algebra is continuous. In that case the inequality (6) has to be satisfied for identical weights  $w_p(\sigma) = w_p(\tau) = w_p(\rho) = w_p$ ,  $p \ge 1$ . If we fix q = 1 then (6) implies  $p \cdot w_{p+1} \le \delta \cdot w_p \cdot w_1$  for  $p \in \mathbb{N}$ . As a consequence we obtain  $w_p \le \delta^{p-1} ((p-1)!)^{-1} w_1, p \ge 1$ . The slowest decrease of the weights which might be possible according to our estimates is therefore  $w_p \sim ((p-1)!)^{-1}$ . The proof that such a solution actually exists follows from the simple estimate  $\binom{m+n}{m} = \frac{(m+n)!}{m!n!} \ge 1$  if  $m, n \ge 0$ . Hence  $(p+q-1)\frac{1}{(p+q-1)!} = \frac{1}{(p+q-2)!} \le \frac{1}{(p-1)!}\frac{1}{(q-1)!}$  is valid for all  $p, q \ge 1$ . Since

$$2^{m+n} \ge \binom{m+n}{m} = \frac{(m+n)!}{m!n!} \ge m+n \text{ if } m, n \ge 1,$$
(8)

also  $(p+q-1)\frac{1}{(p+q)!} < \frac{1}{(p+q-1)!} \le \frac{1}{p!}\frac{1}{q!}$  follows for all  $p, q \ge 1$ . We have therefore derived

**Corollary 2** If the norm is defined with the weights  $w_0 = 1$ ,  $w_n = ((n-1)!)^{-1}$ ,  $n \ge 1$ , or with  $w_0 = 1$ ,  $w_n = (n!)^{-1}$ ,  $n \ge 1$ , the product of the algebra is continuous with the uniform norm estimate

$$||a \circ b|| \le \sqrt{3} ||a|| ||b||.$$
(9)

As a more general class of norms we choose weights

$$w_0 = 1, \ w_n(\sigma, \rho, s) = (n!)^{\sigma} 2^{\rho n} (1+n)^s \text{ if } n \ge 1,$$
 (10)

with real parameters  $\sigma, \rho, s$ . These weights satisfy the inequalities

 $w_n(\sigma_1, \rho_1, s_1) \leq w_n(\sigma_2, \rho_2, s_2)$  if  $\sigma_1 \leq \sigma_2, \rho_1 \leq \rho_2, s_1 \leq s_2$ . We denote by  $||a||_{(\sigma,\rho,s)}$  the norm (4) defined with the weights  $w_n(\sigma, \rho, s)$ . The estimate (8) and the bounds  $\frac{(m+n)!}{m!n!} \geq \frac{(2m)!}{(m!)^2} \geq const \cdot 2^{2m}m^{-\frac{1}{2}}$  if  $n \geq m \geq 1$  and  $1 \leq \frac{(1+m)(1+n)}{1+m+n} \leq 1 + \min(m, n)$  yield inequalities of the type (6) also for these norms. We obtain

$$(p+q-1)w_{p+q}(\sigma,\rho,s) \le \delta w_p(\sigma',\rho',s')w_q(\sigma',\rho',s') \text{ if } p,q \ge 1$$

$$(11)$$

with a constant  $\delta \geq 1$  if  $\sigma = \sigma' = -1$  with  $\rho = \rho' \in \mathbb{R}$  and  $s = s' \leq 0$ , or if  $\sigma = \sigma' < -1$  with  $\rho = \rho' \in \mathbb{R}$  and  $s = s' \in \mathbb{R}$ .

The generalizations of (9) are therefore

$$\|a \circ b\|_{(-1,\rho,s)} \le \sqrt{3} \|a\|_{(-1,\rho,s)} \cdot \|b\|_{(-1,\rho,s)} \text{ if } \rho \in \mathbb{R}, s \le 0,$$
(12)

and

$$\|a \circ b\|_{(\sigma,\rho,s)} \le \gamma \|a\|_{(\sigma,\rho,s)} \cdot \|b\|_{(\sigma,\rho,s)} \text{ if } \sigma < -1, \rho \in \mathbb{R}, s \in \mathbb{R}.$$

$$(13)$$

Here  $\gamma$  takes some value  $\gamma \geq \sqrt{3}$  depending on the choice of the parameters  $\sigma$  and s.

Moreover, the inequalities (11) are valid for  $(\sigma, \rho, s) \neq (\sigma', \rho', s')$  if  $\sigma < \sigma'$  or if  $\sigma = \sigma'$ and  $\rho < \rho'$ . The corresponding estimates for the norms are

$$\|a \circ b\|_{(\sigma,\rho,s)} \le \gamma \|a\|_{(\sigma',\rho',s')} \cdot \|b\|_{(\sigma',\rho',s')} \text{ if } \sigma < \sigma' \text{ for all } \rho, \rho', s, s' \in \mathbb{R},$$
(14)

and

$$\|a \circ b\|_{(\sigma,\rho,s)} \le \gamma \|a\|_{(\sigma,\rho',s')} \cdot \|b\|_{(\sigma,\rho',s')} \quad \text{if } \rho < \rho' \text{ for all } \sigma, s, s' \in \mathbb{R}.$$

$$\tag{15}$$

The value of  $\gamma \geq \sqrt{3}$  depends on the choice of the parameters.

For the tensor algebra and for algebras of symmetrized tensors<sup>5</sup> the Hilbert space  $\mathcal{A}_1 = \mathcal{H}$  generates the whole algebra. Given a (self-adjoint/positive) operator A on  $\mathcal{H}$ , the mapping  $\Gamma(A)e_0 = e_0$  and  $\Gamma(A)(f_1 \circ f_2 \circ \ldots \circ f_n) := (Af_1) \circ (Af_2) \circ \ldots \circ (Af_n)$  for  $f_{\mu} \in \mathcal{H}, \mu = 1, \ldots, n$ , and  $n \in \mathbb{N}$ , defines a unique (self-adjoint/positive) operator  $\Gamma(A)$  on the algebra  $\mathcal{A}$ , which satisfies the relation

$$\Gamma(A)(a \circ b) = (\Gamma(A)a) \circ (\Gamma(A)b).$$
(16)

The norms (4) with the weights (10) are then easily generalized to

$$\|a\|_{(\sigma,\rho,s)}^{2} = \sum_{n=0}^{\infty} (n!)^{\sigma} \|(\Gamma(A))^{\rho} a_{n}\|_{n}^{2} (1+n)^{s} \text{ if } a = \sum_{n=0}^{\infty} a_{n}, \ a_{n} \in \mathcal{A}_{n}.$$
(17)

If A is an invertible positive operator with lower bound  $A \ge 2 \cdot id$ , then  $\Gamma(A)$  satisfies  $\|(\Gamma(A))^{-\rho} a\|_n \le 2^{-n\rho} \|a\|_n$  for  $a \in \mathcal{A}_n$  if  $\rho \ge 0$ . This bound and the relation (16) imply that the estimates (12),(13) and (15) are also valid for the norms (17), moreover (14) holds if  $\rho \le \rho'$ .

If  $A^{-1}$  is a Hilbert-Schmidt operator then a family of norms (17) can be used to define a nuclear topology on the algebra  $\mathcal{A}$ . For the symmetric tensor algebra that has been done in the white noise calculus and in the Malliavin calculus, see e.g. [1] [11] [20]. For the algebra of antisymmetric tensors and for the superalgebras such nuclear topologies can be found in [12] and in [6]. But the estimates of these references are not strong enough to derive the results with a single Hilbert norm as presented in Corollary 2 and in eqs. (12) and (13).

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<sup>&</sup>lt;sup>5</sup>This class of algebras does not only include the algebra of symmetric tensors and the algebra of antisymmetric tensors (Grassmann algebra), but also the  $\mathbb{Z}_2$ -graded algebras (superalgebras) used in supersymmetric quantum field theory.

## References

- A. Arai and I. Mitoma. Comparison and nuclearity of spaces of differential forms on topological vector spaces. J. Funct. Anal., 111:278-294, 1993.
- [2] J. C. Baez, I. E. Segal, and Z. Zhou. Introduction to Algebraic and Constructive Quantum Field Theory. Princeton University Press, Princeton, 1992.
- [3] F. A. Berezin. The Method of Second Quantization. Academic Press, New York, 1966.
- [4] F. A. Berezin. Introduction to Superanalysis. Reidel, Dordrecht, 1987.
- [5] B. DeWitt. Supermanifolds. CUP, Cambridge, 1984.
- [6] Z. Haba and J. Kupsch. Supersymmetry in euclidean quantum field theory. Fortschr. Phys., 43:41-66, 1995.
- [7] T. Hida, H.-H. Kuo, J. Potthoff, and L. Streit. White Noise. Kluwer, Dordrecht, 1993.
- [8] A. Jadczyk and K. Pilch. Superspaces and supersymmetries. Commun. Math. Phys., 78:373-390, 1981.
- [9] Y. Le Jan. On the Fock space representation of functionals of the occupation number field and their renormalization. J. Funct. Anal., 80:88–108, 1988.
- [10] A. Yu. Khrennikov. Functional superanalysis. Russian Math. Surveys, 43:103–137, 1988.
- [11] Yu G. Kondratev and L. Streit. Spaces of white noise distributions: constructions, descriptions, applications. I. *Rep. Math. Phys.*, 33:341–366, 1993.
- [12] P. Krée. Méthodes fonctionelles en analyse de dimension infinie et holomorphie anticommutative. In Séminaire P. Lelong et H. Skoda (Analyse) Année 1976/77, Lect. Notes Math. 694, pages 134-171. Springer, Berlin, 1978.
- [13] J. Kupsch. A probabilistic formulation of bosonic and fermionic integration. Rev. Math. Phys., 2:457-477, 1990.
- [14] J. L. Martin. Generalized classical dynamics, and the "classical analogue" of a Fermi oscillator. Proc. Roy. Soc. (London), A251:536-542, 1959.
- [15] P. A. Meyer. Quantum Probability for Probabilists. Lect. Notes in Math. 1538. Springer, Berlin, 1993.
- [16] A. Rogers. A global theory of supermanifolds. J. Math. Phys., 21:1352–1365, 1980.
- [17] A. Rogers. Graded manifolds, supermanifolds and infinite-dimensional Grassmann algebras. Commun. Math. Phys., 105:375–384, 1986.

- [18] O. G. Smolyanov and E. T. Shavgulidze. The Fourier transform and pseudodifferential operators in superanalysis. Soviet Math. Dokl., 37:476–481, 1988.
- [19] V. S. Vladimirov and I. V. Volovich. On the definition of the integral in superspace. Soviet Math. Dokl., 32:817–819, 1985.
- [20] S. Watanabe. Lectures on Stochastic Differential Equations and Malliavin Calculus. Tata Institute, Bombay, and Springer-Verlag, Berlin, 1984.